

# Universal Aging Dynamics in $\text{SiO}_2$

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## Outline:

1. Microscopic: Single Particle Jump Dynamics
2. Scaling:
  - ▶ global incoherent intermediate scattering function,
  - ▶ dynamic susceptibility
  - ▶ local incoherent intermediate scattering function

## Model: BKS Potential

[B.W.H. van Beest *et al.*, PRL 64, 1955 (1990)]

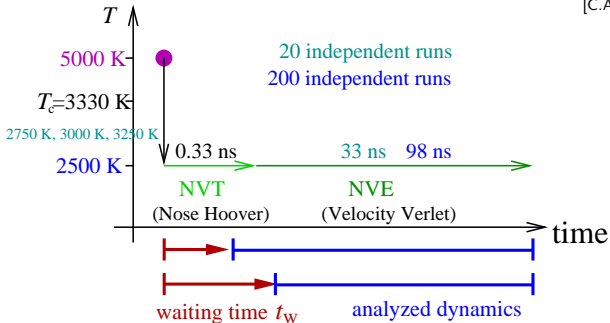
$$\phi(r_{ij}) = \frac{q_i q_j e^2}{r_{ij}} + A_{ij} e^{-B_{ij} r_{ij}} - \frac{C_{ij}}{r_{ij}^6}$$

112 Si & 224 O

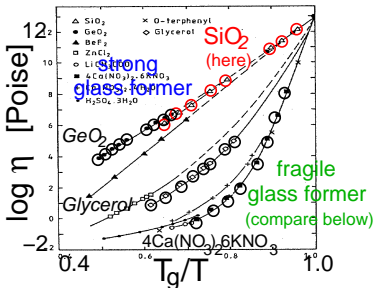
$\rho = 2.32 \text{ g/cm}^3$

$T_c = 3330 \text{ K}$

## Simulation Runs:

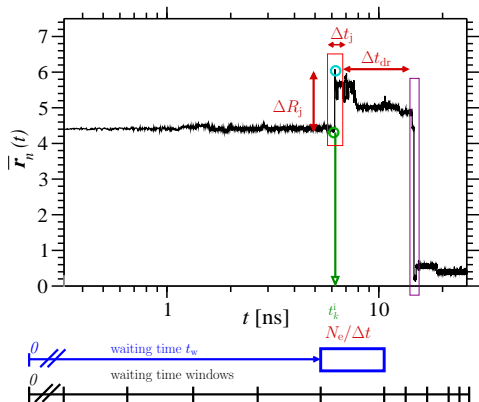
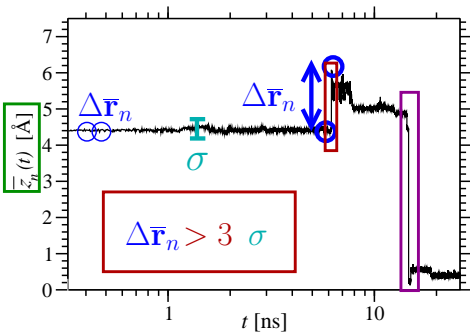


## Dynamics:



[C.A. Angell *et al.* 1976]

# 1. Microscopic: Jump Definition & Jump Statistics



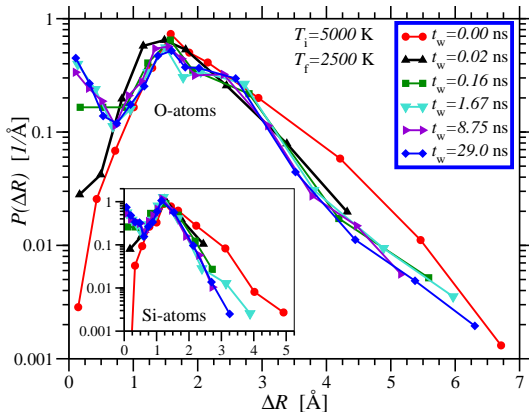
[KVL, R. Bjorkquist, L.M. Chambers, PRL 110, 017801 (2013)]

see also [KVL, J. Chem. Phys. 121, 4781 (2004)]

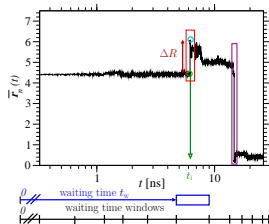
jumps  $\longleftrightarrow$  defects [KVL & A. Zippelius, PRE 88, 052145 (2013)]

# Jump Length Distribution

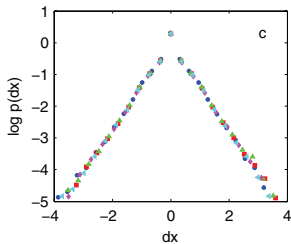
strong glass former  $\text{SiO}_2$ :



►  $P(\Delta R)$  independent of  $t_w$

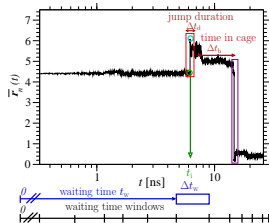
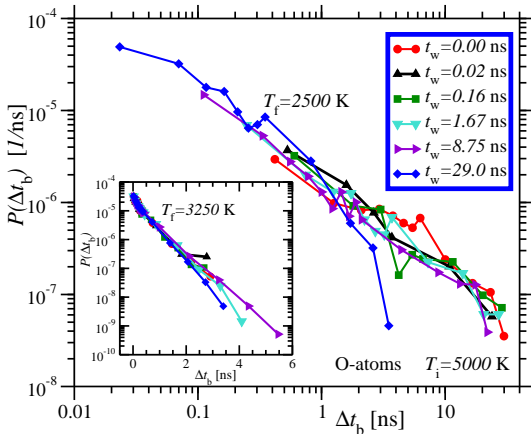


► compare fragile glassformer binary LJ (& polymer) [Warren & Rottler, EPL(2009)]

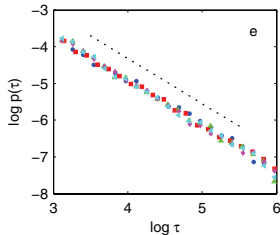


# Distribution of Time in Cage $P(\Delta t_b)$ : $t_w$ varied

strong glass former  $\text{SiO}_2$ :



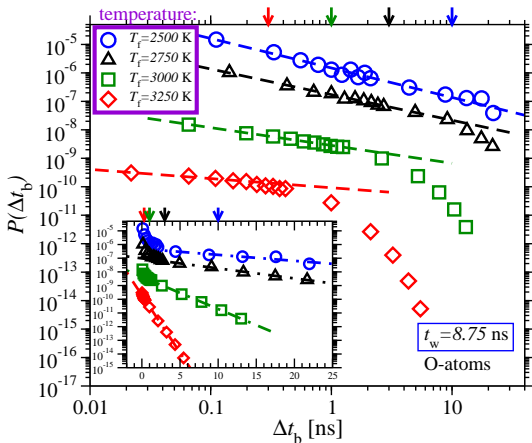
► compare fragile glassformer (binary LJ &) polymer [Warren & Rottler, EPL(2009)]



►  $P(\Delta t_b)$  independent of  $t_w$ !

# Distribution of Time in Cage $P(\Delta t_b)$ : $T_f$ varied

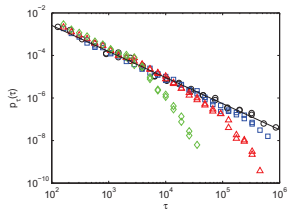
strong glass former  $\text{SiO}_2$ :



► crossover

- power law to exponential
- at  $t_{\text{cross}} \approx t_{\text{eq}}^C$

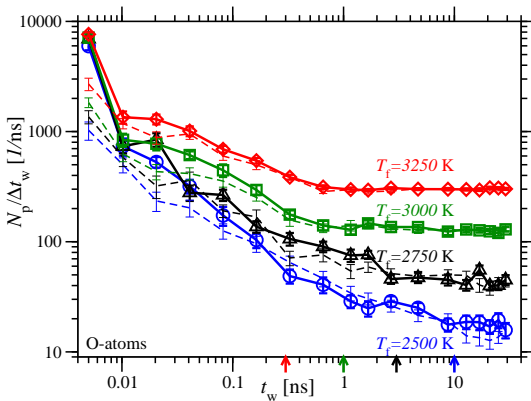
► compare fragile glassformer binary LJ [Warren & Rottler, '13]



polymer; CTRW

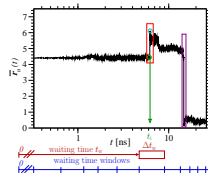
[Helfferrich et al. '15]

# Number of Jumping Particles per Time



→ Only  $t_w$ -dependence:  $N_p / \Delta t_w$   
(not  $P(\Delta R)$  and  $P(\Delta t_b)$ )

[KVL, R. Bjorkquist, L.M. Chambers, PRL 110, 017801 (2013)]



►  $N_p / \Delta t_w$  depends strongly on waiting time  $t_w$

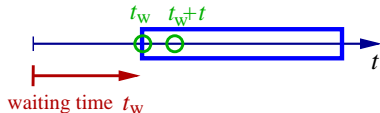
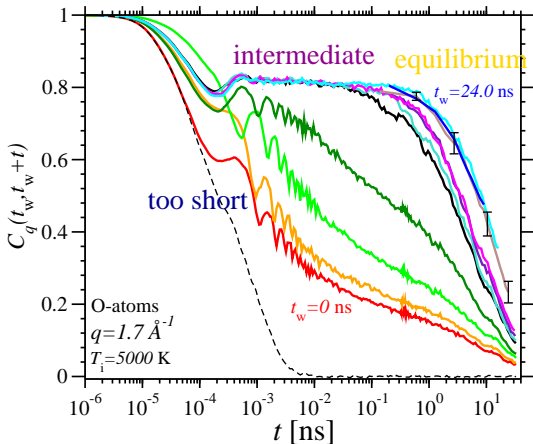
- $N_p / \Delta t_w$  decreasing with increasing  $t_w$
- compare: soft colloids [Yunker et al., PRL (2009)]
- compare polymer; CTRW [Helfferich et al., EPL 2015]

► equilibration at  $t_{eq}^j$   
 $t_{eq}^j \approx t_{eq}^C$  (arrows)

## 2. Scaling: Global Incoherent Intermediate Scattering Fct.

$$f_s^\alpha(t_w, t_w + t, \mathbf{q}) = \frac{1}{N_\alpha} \sum_{j=1}^{N_\alpha} \cos \{ \mathbf{q} \cdot (\mathbf{r}_j(t_w + t) - \mathbf{r}_j(t_w)) \}$$

$$C^\alpha = \langle f_s^\alpha(t_w, t_w + t, \mathbf{q}) \rangle$$



- ▶  $C_q(t_w, t_w + t)$  depends on waiting time  $t_w$  (colors)
- ▶ three time windows:
  - ▶ too short
  - ▶ intermediate (scaling)
  - ▶ equilibrium ( $t_{eq}^C$ )
- ▶ equilibrium curve included in scaling

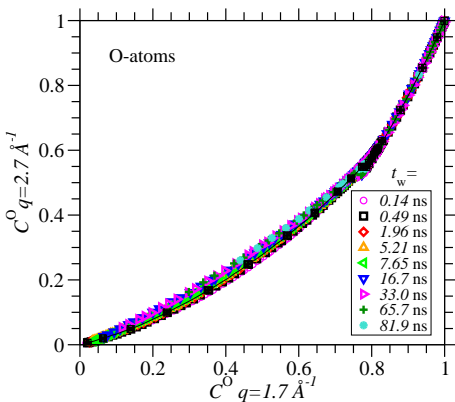


## 2. Scaling: Global Incoherent Intermediate Scattering Fct.

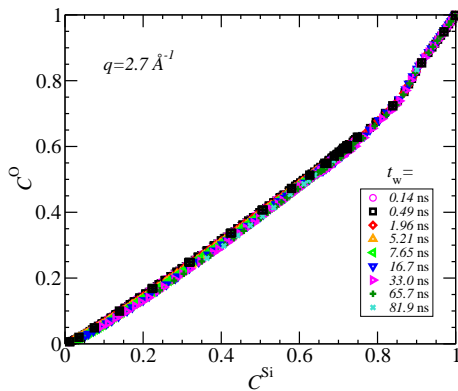
$$f_s^\alpha(t_w, t_w + t, \mathbf{q}) = \frac{1}{N_\alpha} \sum_{j=1}^{N_\alpha} \cos \{ \mathbf{q} \cdot (\mathbf{r}_j(t_w + t) - \mathbf{r}_j(t_w)) \}$$

$$C^\alpha = \langle f_s^\alpha(t_w, t_w + t, \mathbf{q}) \rangle$$

$\alpha =$  particle type Si,O



$$C^\alpha = C^\alpha(z(t_w, t), q)$$



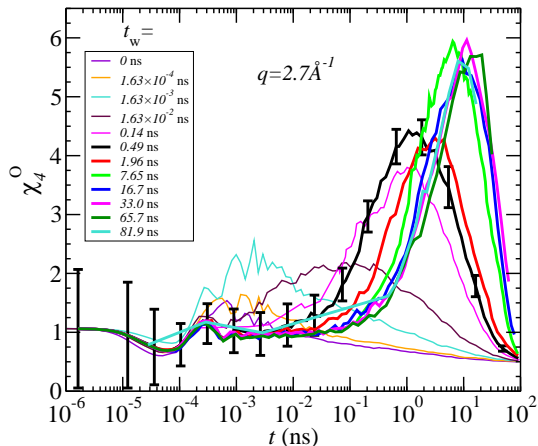
$$C_q = C_q(z(t_w, t), \alpha)$$

$$C = C(z(t_w, t), q, \alpha)$$

# Dynamic Susceptibility

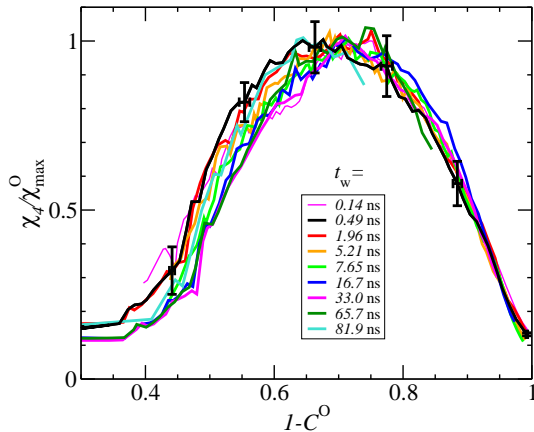
$$f_s^\alpha(t_w, t_w + t, \mathbf{q}) = \frac{1}{N_\alpha} \sum_{j=1}^{N_\alpha} \cos \{ \mathbf{q} \cdot (\mathbf{r}_j(t_w + t) - \mathbf{r}_j(t_w)) \}$$

$$\chi_4^\alpha(t_w, t_w + t, q) = N_\alpha \left[ \langle (f_s^\alpha)^2 \rangle - (\langle f_s^\alpha \rangle)^2 \right]$$



# Scaling of Dynamic Susceptibility

$$\chi_4^\alpha(t_w, t_w + t, q) = N_\alpha \left[ \langle (f_s^\alpha)^2 \rangle - (\langle f_s^\alpha \rangle)^2 \right]$$



► theory:

$$\chi_4/\chi_4^{\max} = \chi_4/\chi_4^{\max}(C)$$

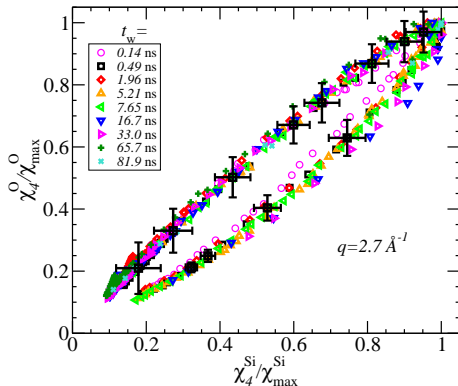
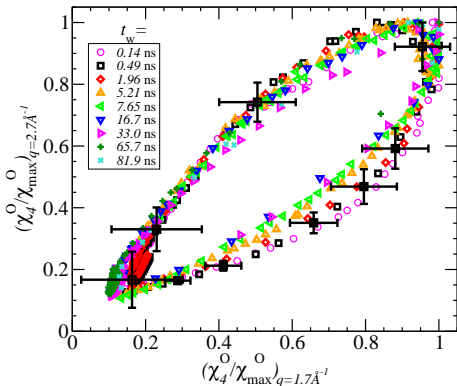
[Castillo et.al '02 – '04,'08]

► compare fragile glassformers:

[Parsaeian & Castillo, '08,'09]

# Scaling of Dynamic Susceptibility

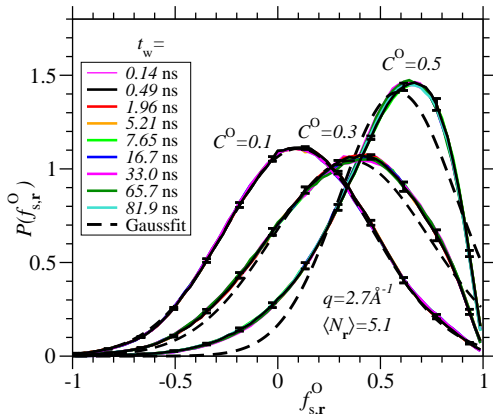
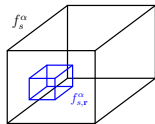
$$\chi_4^\alpha(t_w, t_w + t, q) = N_\alpha \left[ \langle (f_s^\alpha)^2 \rangle - (\langle f_s^\alpha \rangle)^2 \right]$$



$$\longrightarrow \chi_4 / \chi_4^{\max} = \chi_4 / \chi_4^{\max} (C(z(t_w, t), q, \alpha))$$

# Local Incoherent Intermediate Scattering Function

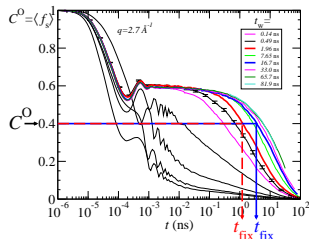
$$f_{s,r}^\alpha(t_w, t_w + t, \mathbf{q}) = \frac{1}{N_r^\alpha} \sum_{\mathbf{r}_j(t_w) \in B_r} \cos(\mathbf{q} \cdot [\mathbf{r}_j(t_w + t) - \mathbf{r}_j(t_w)])$$



Theory:

$P(f_{s,r})$  scales with  $C^\alpha$

[Castillo et al.'02,'03; Chamon et al. '04]



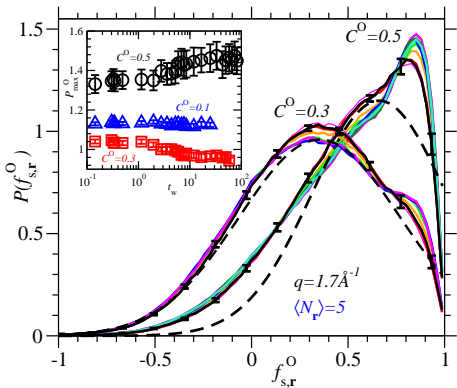
compare fragile & spin glasses:

[Castillo & Parsaeian, '07, '09] & [Castillo &

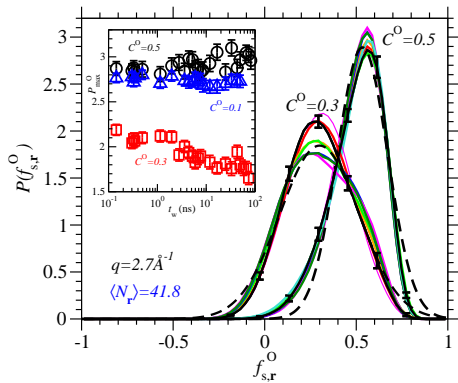
Chamon et al., '02 – '04]

# NOT Scaling of Local Incoh. Intermediate Scattering Fct.

$$f_{s,r}^\alpha(t_w, t_w + t, \mathbf{q}) = \frac{1}{N_r^\alpha} \sum_{\mathbf{r}_j(t_w) \in B_r} \cos(\mathbf{q} \cdot [\mathbf{r}_j(t_w + t) - \mathbf{r}_j(t_w)])$$



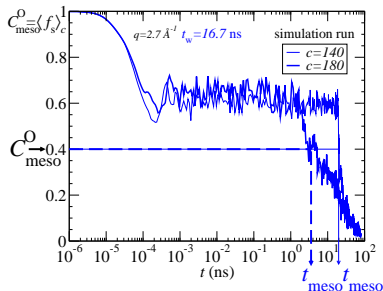
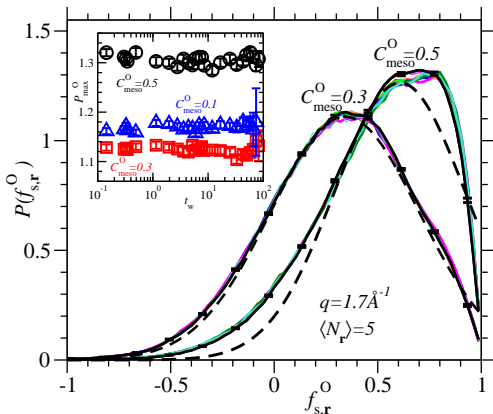
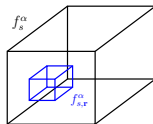
wave vector  $q$  too large



subbox  $B_r$  too large

# Scaling of Local Incoh. Intermediate Scattering Fct.

$$f_{s,r}^\alpha(t_w, t_w + t, \mathbf{q}) = \frac{1}{N_r^\alpha} \sum_{\mathbf{r}_j(t_w) \in B_r} \cos(\mathbf{q} \cdot [\mathbf{r}_j(t_w + t) - \mathbf{r}_j(t_w)])$$



→  $P(f_{s,r}^\alpha(t_w, t))$  scales with  $C_{\text{meso}}^\alpha(t_w, t)$

## Summary:

Aging Dynamics of Strong Glass Former SiO<sub>2</sub>:

1. Microscopic: Single Particle Jump Dynamics
  - ▶ Only  $t_w$ -dependence:  $N_p/\Delta t_w$  (not  $P(\Delta R)$  and  $P(\Delta t_b)$ )
2. Scaling:
  - ▶ global incoh. interm. scattering function  $C = C(z(t_w, t), q, \alpha)$   
 $\alpha = \text{Si, O}$
  - ▶ dynamic susceptibility  $\chi_4/\chi_4^{\max} = \chi_4/\chi_4^{\max}(C(z(t_w, t), q, \alpha))$
  - ▶ local incoh. interm. sc. fct. distribution  $P(f_{s,r})$  scales with  $C$
3. similar aging dynamics of strong & fragile glasses

## Acknowledgments:

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